

Coronal heating via nanoflares:

Spontaneous Current Sheets Unavoidable in 3D fields!

Åse Marit Janse

Advanced Study Program / High Altitude Observatory

National Center for Atmospheric Research



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Coronal heating

- Energy flux needed
 - ~ 10^4 W m⁻² in active regions
 - ~ 10^2 W m⁻² in the quiet sun (Withbroe & Noyes 1977)
- Photospheric motions moving the footpoints of coronal magnetic fields increases the magnetic free energy by
 - Poynting flux: ~10⁴ W m⁻²
- Problem: How is this energy dissipated in the corona?
 Nanoflares?

Nanoflares: Dissipation of current sheets

- CS Formation:
 - Each individual flux tube moves independently of its neighbours (due to the photospheric motions)
 - In the corona where these flux tubes expand against each other
 - In general, the magnetic fields at the boundary are not aligned
 - \Rightarrow Tangential discontinuities (i.e. current sheets)
- CS Dissipation:
 - Release (built up) magnetic free energy

If nanoflares is the solution then...

- First of all
 - Current sheets form easily!
 - Their formation is spatially extensive
- Furthermore
 - They form & dissipate quickly enough to release the desired energy flux

Spontaneous current sheets

In the highly electrically conducting corona: -The magnetic field is frozen into the plasma -Preservation of field topology (footpoint map and twist)

Parker theory: A field of fixed field topology and B_n can be in a continuous state in one equilibrium but may have to contain inevitable tangential discontinuities on transition to another equilibrium

Field deformation: Footpoint motion, volume change... \Rightarrow Current sheet formation is unavoidable!



In 3D fields current sheets form easily Example: Volume change

(Janse & Low2009, ApJ 690)

A potential field inside a cylinder of perfectly conducting fluid. The field is anchored at the cylinder ends: Topology invariant to a change in L (Field remains untwisted & footpoint map preserved)

Fully 3D field, i.e. azimuthal dependency No symmetries!





In 3D fields current sheets form easily Example: Volume change

A potential field inside a cylinder of perfectly conducting fluid. The field is anchored at the cylinder ends: Topology invariant to a change in L (Field remains untwisted & footpoint map preserved)

Then $L_1 \rightarrow L_2$: The deform field seeks a new equilibrium

Can **B**_{deformed} evolve, under the frozen-in condition, into **B**_{potential}?

(**B**_{potential}: the only <u>continuous</u> untwisted equilibrium state in the deformed volume)



Volume change in 3D fields: Current sheets must form and dissipate!

z=L



Implications:

- 1. Under the frozen-in cond: **B_{deformed}** <u>cannot</u> reach **B**_{potential}
- 2. We suggest that current sheets must form throughout the field, whose dissipation can then change T(B_{deformed}) to match T(B_{potential})



Volume change in 2D fields

An <u>axisymmetric</u> potential field inside a cylinder Topology invariant to a change of L

Then $L_1 \rightarrow L_2$: Can $\mathbf{B}_{deformed}$ evolve, under the frozen-in condition, into $\mathbf{B}_{potential}$?

If no neutral points: $T(B_{deformed})=T(B_{potential})$, T fixed by B_n by No current sheets



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If neutral points:

 $T(B_{deformed}) \neq T(B_{potential}),$ Current sheets only at specific locations: in the vicinity of <u>separatrix surfaces</u>

Hard X-ray flare footpoints & current sheet formation



Hannah et al. 2008, A&A 481

Footpoint brightening of macroscopic size + electrons channeled along the field lines ⇒ suggests spatially extensive current sheet formation (for 3D twisted fields)

Concluding remarks ...

Current sheet formation in 3D fields: – Current sheets form readily

Current sheets form throughout the field

Is the dissipated energy large enough?